

The Classification of Topological Orders


Math Dept Colloquium
The Ohio State University
1 April 2021

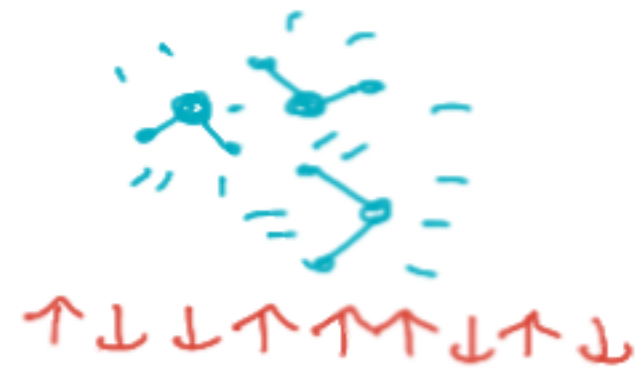
Theo Johnson-Freyd
Dalhousie University &
Perimeter Institute

These slides available at <http://categorified.net/OSU.pdf>

Phases of matter

Microscopically, a system of matter might be described by

- molecules bumping into each other
- spins at each site in a lattice
- quantum fields 
- ...



"The UV"

The phase of the system is its long-range effective macroscopic description.

"The IR"

Solid or fluid? Ferromagnetic? Conductive?

Phases of matter

A humongous amount of physics research focuses on two overlapping questions:

★ Classify phases of matter

★ Understand the correspondence between UV and IR descriptions.

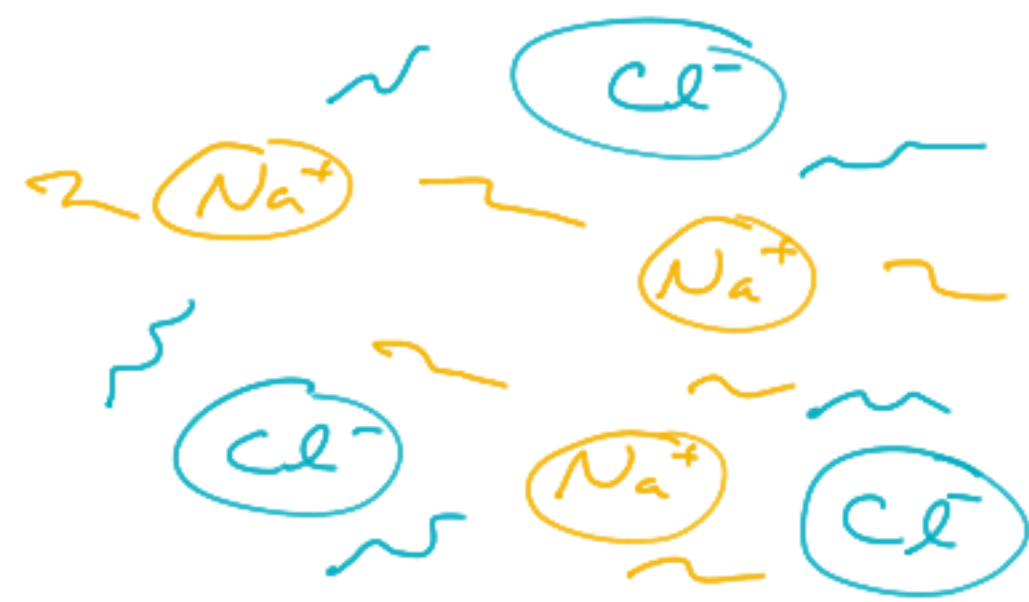
↳ In good situations, small changes to UV parameters do not affect the IR description.

So the UV-to-IR map is many-to-one.

↳ The UV-to-IR map is not directly computable with perturbation theory, the usual tool in QFT.

Landau Paradigm

Consider anhydrous NaCl. Microscopically, this is a system of ions bumping into each other.



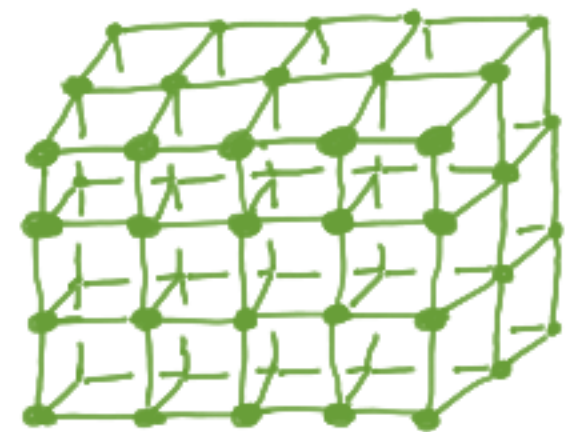
At high temperatures, it's macroscopically a liquid.



In particular, its phase is symmetric under all of $O(3) \times \mathbb{R}^3$.

rotations + reflections translations

At low temperatures, it's macroscopically a crystal.



This is still effectively translation-invariant.

But the $O(3)$ symmetry spontaneously breaks to $\text{Aut}(\boxplus) \cong C_2 \times S_4$
↳ measure w/ opacity, direction of cracks, etc.

same finite gp of order 48.

The crystal phase is ordered.

Landau Paradigm

The **Landau Paradigm** says that phases are classified by their patterns of **symmetry** and **symmetry breaking**.

It is **remarkably effective**, able to explain

- solid/liquid/gas
- (anti)ferromagnetism
- conductor/insulator
- relative strength of electro/weak forces

It suggests a solution to the UV-to-IR question:

(1) Identify all symmetries of the UV description.

Work out basic properties like whether there is an **anomaly**:
is the action on the Hilbert space linear or projective?

(2) Identify all possible patterns of symmetry breaking.

Describe the phase with that **order**.

Topologically Ordered Phases

Phases that violate the Landau Paradigm emerged around the turn of the century. These phases are deeply quantum, with long range entanglement.

Indeed, all of the physical data is encoded in the long range entanglement, and so the phase is locally trivial. In particular, there are

no ^{nontriv} local operators in the IR description.

The local physics is invariant under all diffeomorphisms: it is a topological liquid.

Physics defn: A liquid is topological when its stress-energy tensor is central, i.e. $\propto \langle \text{vac} \rangle$.

"Ordered by topology rather than symmetry"

← useful for quantum computing!

The local UV operators are "gapped out."

The energy penalty required is too high for IR physics.

Topologically Ordered Phases

How to tell that such a phase is nontrivial?

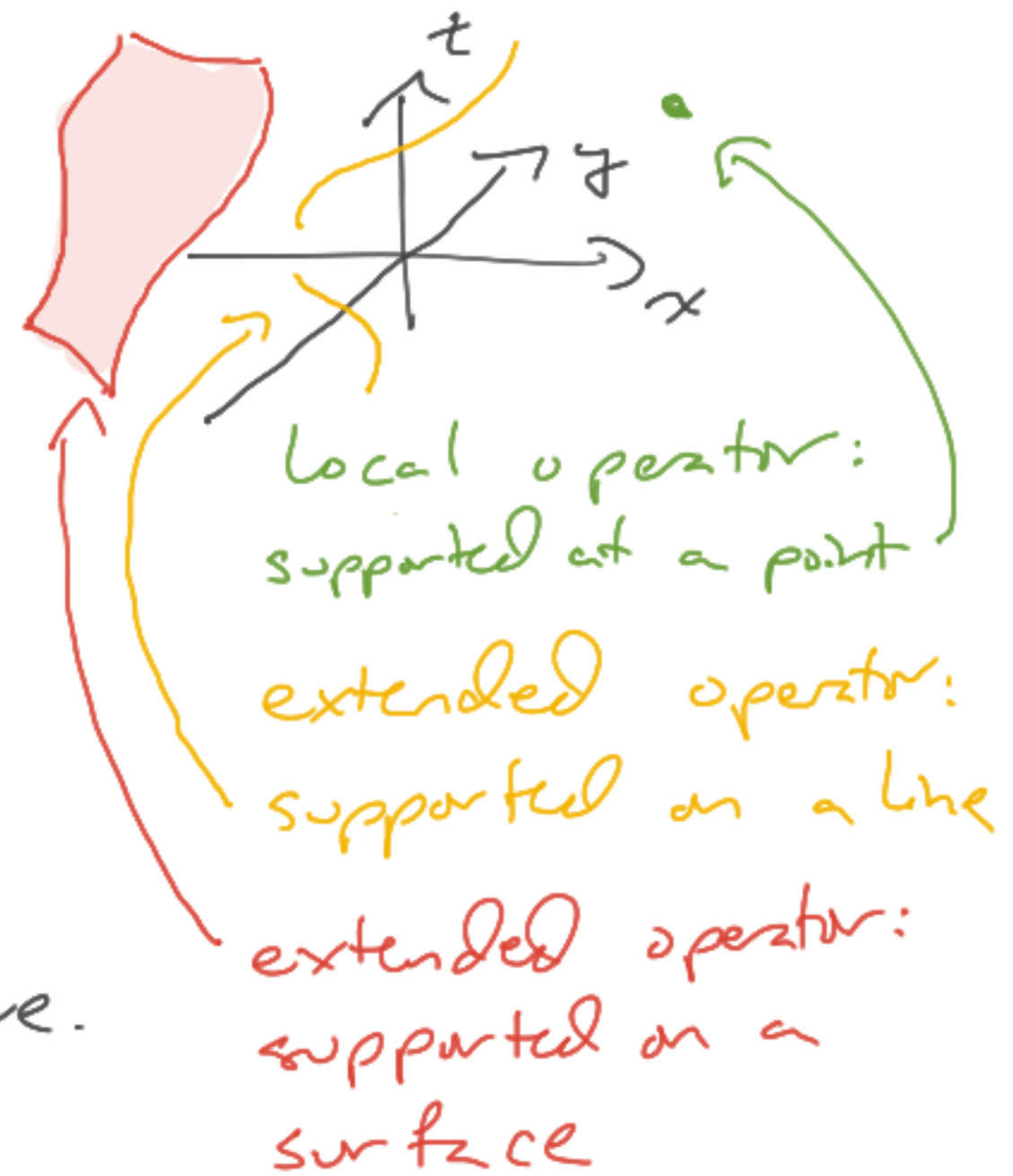
It has **extended operators**.

Non topological example: In pure Maxwell theory (i.e. U(1) gauge theory) in d dimensions, there are **test electrons** which measure the holonomy of the photon field along a curve.

Physically, inserting this operator measures the response an **electrically charged particle** would feel if it underwent that path in spacetime.

There are also **test magnetons**, which are $(d-3)$ -dim (magnetic monopole)

These operators do not commute: they detect each other.



In Maxwell theory, these operators are \oint local operator.

Topologically Ordered Phases

How to tell that such a phase is nontrivial?

It has **extended operators**.

Topological example: In the Toric Code

(i.e. \mathbb{Z}_2 gauge theory) in 2 dimensions

there are particle (aka line) **test electrons**

and $(d-2)$ -dimensional **test magnetons**.

$d=2+1$ commutation relation:

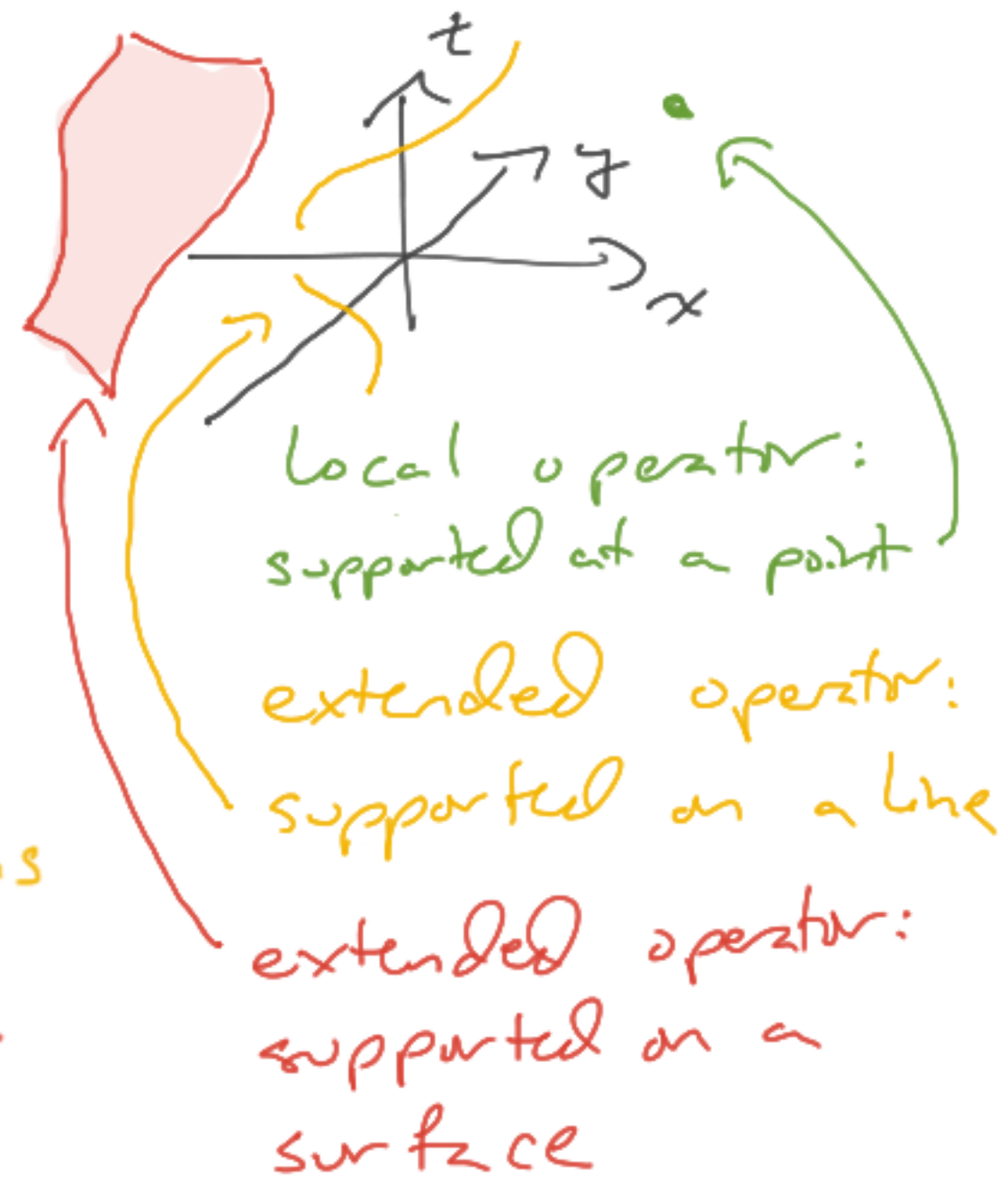
 = -

analogue of " $xy = -yx$ "

 = + ,  = +

analogues of " $x^2=1$ ".

Other coefficients are also allowed,
e.g. in "twisted gauge theory".



Mathematical Axioms

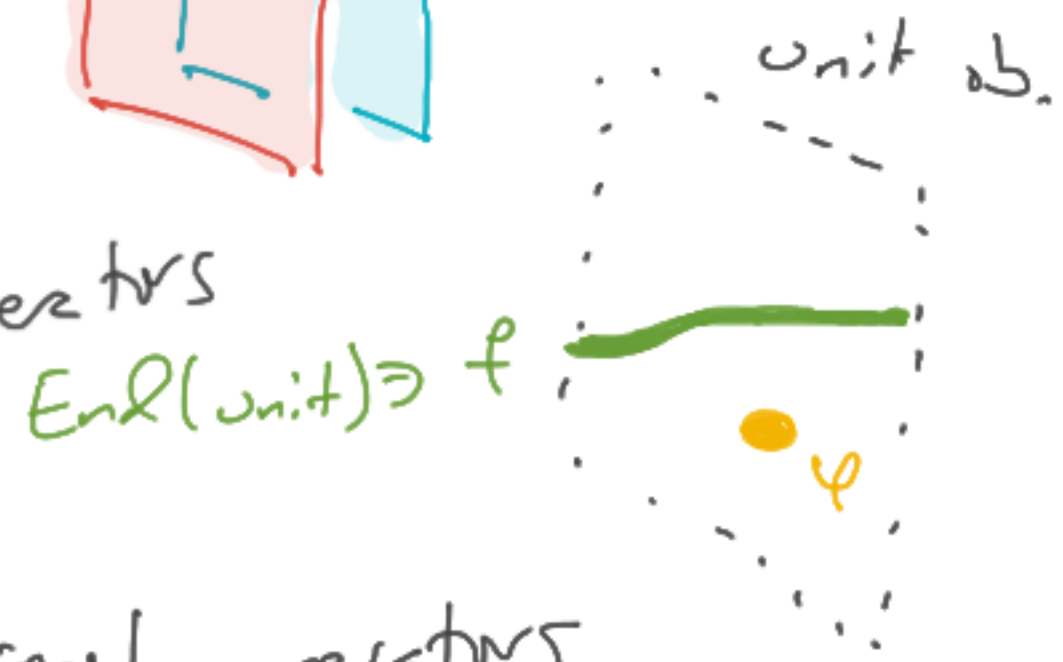
aka a topological order

A topological quantum liquid phase is determined by its "algebra" of extended operators. In $(n+1)$ dimensions, this "algebra" is a monoidal n -category \mathcal{A} .

- $\text{Ob}(\mathcal{A}) = n$ -dimensional operators



- $\text{Ob}(\text{End}_{\mathcal{A}}(\text{unit})) = (n-1)$ -dimensional operators



- $\text{Ob}(\text{End}_{\text{End}_{\mathcal{A}}(\text{unit})}(\text{unit})) = (n-2)$ -dimensional operators

$$\psi \in \text{End}_{\text{End}(\text{unit})}(\text{unit})$$

- Etc.

Mathematical Axioms

This monoidal n -category A should be

Quantum: Linear, additive, and Karoubi complete

↙ superpositions are allowed

Topological: Very strong finiteness conditions

↪ i.e. rigid finite semisimple, i.e. "multifusion"

i.e. central

Remotely detectable: The only invisible operators are scalar multiples of identity.

Robust: The only point operators are scalar multiples of identity

↪ otherwise, the system is in a critical state between two phases.

← forced by asking to be able to place phase on spacetimes like



← "fusion"

A is a higher categorical "central simple algebra".

← "Heisenberg uncertainty"

Classification

$(0+1)D: \{*\}$. Only point operators. Robustness \Rightarrow all trivial

$(1+1)D: \{*\}$. Point and line operators.
Remote detectability \Rightarrow line operators separate points
No points. So no lines.

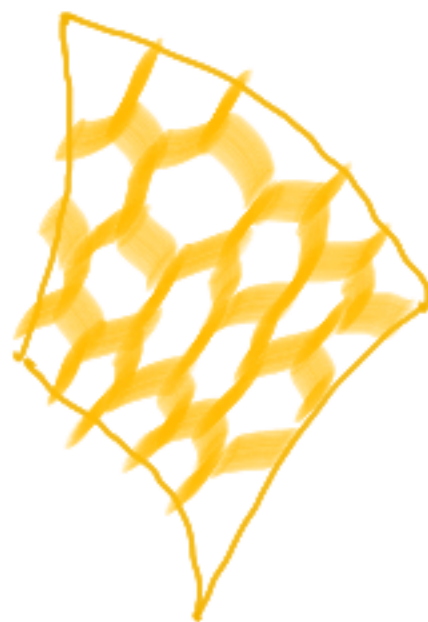
In $(0+1)D$ and $(1+1)D$, there is
no inherent topological order.

But symmetry-enriched systems are interesting because there can be anomalies.

$(2+1)D: \{MTCs\}$.
[wen]

Robust + remote detect \Rightarrow
all surface operators are
networks of line operators

But the line operators
are very rich.



- ∞ many.
- classification seems completely wild.

Classification

(3+1)D:

[Lan, Kong, Wen]

[Lan, Wen]

[JF]

- 3D operators are networks of 2D operators. so all data in the lines and surfaces.
- Line operators are symmetric monoidal.

↳ so the bosonic lines are $\text{Rep}(G)$ for some finite gp G .

"charged test bosons"

↳ can trigger a phase transition which "ungauges" a G -action.

in a G -gauge theory.

Original top order with $\text{Rep}(G)$ lines



almost trivial system w/ G -sym.

↑ phase transition interface, a gapped interface.

Then: If no fermionic lines, then ordinary gauge theory classified by $H^4(G; \mathbb{U}(1))$.

"super cohomology"

Then [JF - Reutter]: These two choices cannot be separated by a gapped interface. →

- If fermionic lines, then two choices for ungauged system, and $SH^4(G)$ choices for gauge theory action.

↙ A certain generalized coh. th.

Classification

(4+1)D: After coupling to a spin structure, "local fermion"
[JF- γ_0] all phases built by
 \hookrightarrow gauging a "1-form symmetry"
 \hookrightarrow then gauging a "0-form symmetry"

The 0-form sym may act by electromagnetic duality on the 1-form gauge theory.

There is some manageable redundancy in this classification.

Before coupling to a spin structure, there are infinitely many phases which cannot be related by a gapped interface / phase transition.

Related to 2-torsion in the fermionic Witt group.

Higher Categorical Landau Paradigm

Symmetries can be encoded by codimension-1 symmetry defects.

They are invertible and topological.

Higher homotopical symmetry

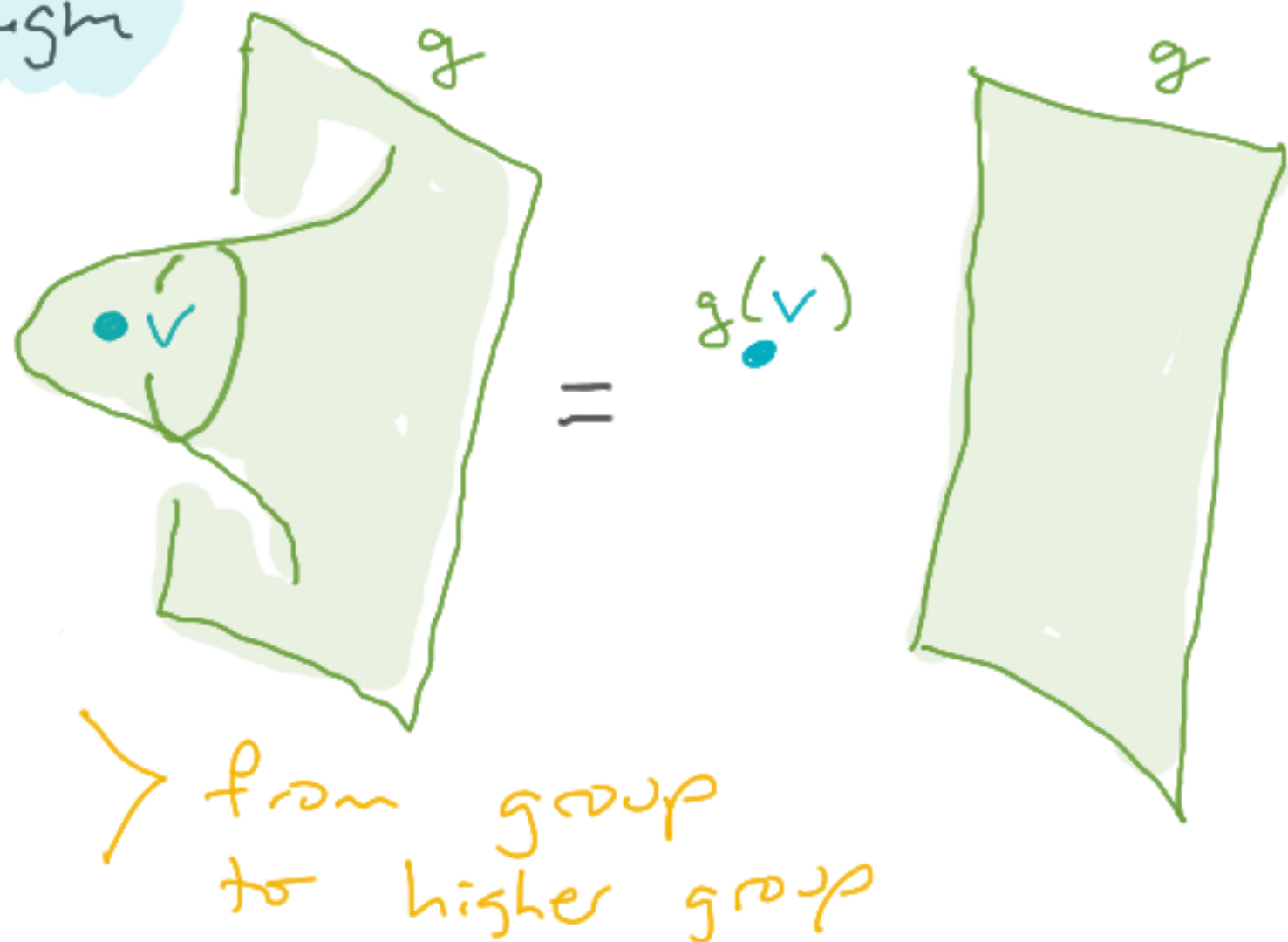
(\rightsquigarrow) higher codimensional sym. defects.

Quantum superposition (\rightsquigarrow) $g \cdot h = x \oplus y \oplus \dots$

(\rightsquigarrow) non-invertible "symmetries"

Quantum field theory requires both generalizations.

"Categorical symmetry" restores the Landau Paradigm.



from group to higher group

from group to algebra.

higher algebra aka n -cat!