

Global Categorical Symmetries

PSI Advertisement, 19 Sept 2022

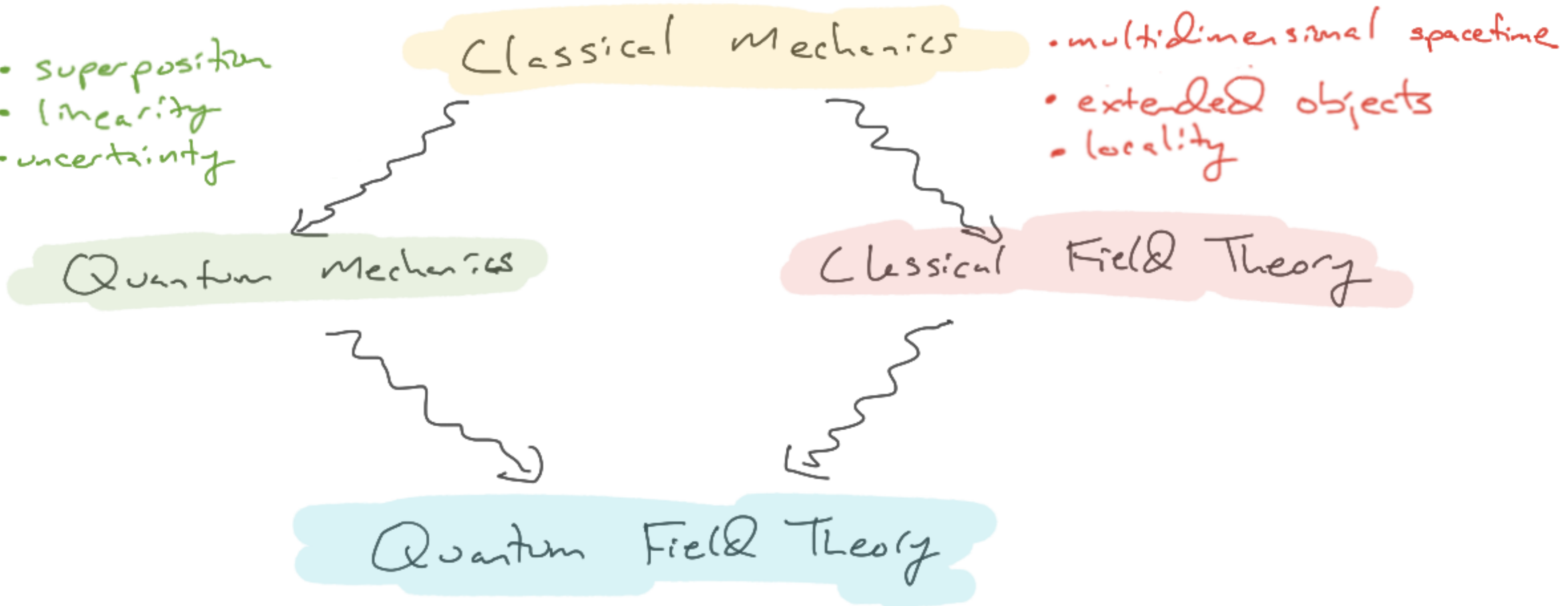
Theo Johnson-Freyd theo.jf@pitp.ca

If you want to chat further, I'll be
in town Oct 4-7. Or send an email.

These slides: categorified.net/PSI-ad.pdf

Symmetries are present in all subdisciplines of math + physics, and essentially every person at PI uses them in their research. I want to tell you about a kind of symmetry that is fundamentally quantum field theoretical. Recall that QFT generalizes classical mechanics in two orthogonal directions:

- superposition
- linearity
- uncertainty



Classical symmetry operations are invertible, and the "temporal" order of composition matters. The abstract compositional structure of classical symmetries is controlled by a group. QFT provides two orthogonal generalizations:

- linearity
- noninvertibility

Noncommutative Algebra

Group Theory

- multidimensional composition
- extended objects
- locality

Homotopy Theory

Category Theory

The word "categorical" in GCS is doing double duty, referring to both of these generalizations at once.

The word "global" also does double duty. On one hand, it is a synonym for "flavour", as opposed to gauged / colour symmetries. But it also emphasizes an interest in the global structure of the symmetry "group". Ever since Noether, physicists have focused on infinitesimal symmetries, controlled by Lie algebras and implemented by Noether currents. Lie algebras do not see global structure, e.g. $SO(n)$ vs. $O(n)$ vs. $Spin(n)$.

I focus on an extreme case of this:
finite symmetries, with no* infinitesimal structure.

* In characteristic zero. An amazing feature of finite groups is that a finite group has inf' structure in char dividing its order.

Even though finite symmetries lack Noether currents, they still have a Noether Theorem, because the "charge" operators

$$Q(\Sigma^2) = \int_{\Sigma} \star J$$

← Noether current
← Hodge star

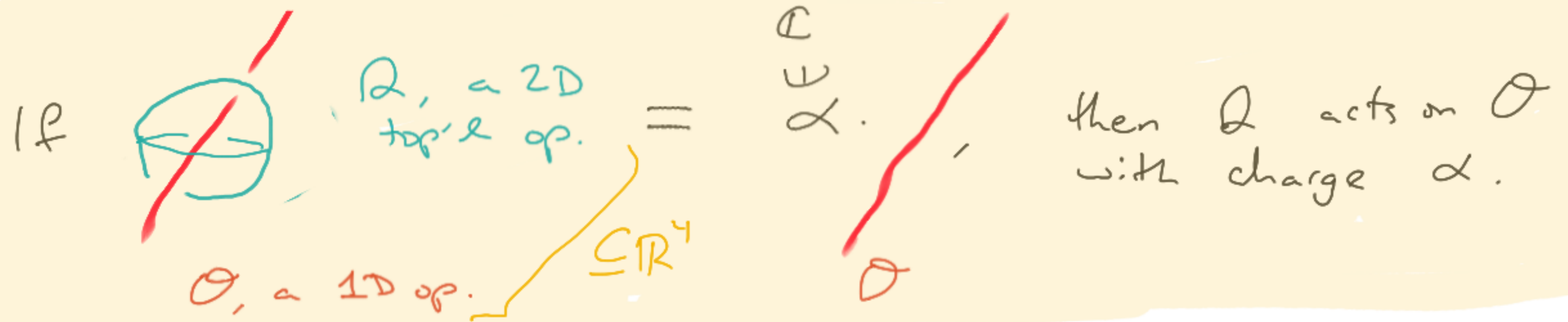
still make sense. Noether: $d \star J = 0$, so $[Q, T^{\mu\nu}] = 0$.

$$\lim_{\ell \rightarrow 0} \frac{1}{\ell} \left(\int_{\Sigma} \star J \Big|_{\ell} - \int_{\Sigma} \star J \Big|_{\ell} \right) = \int_{\Sigma} \star J$$

In other words, Q is topological.

Defn: A p-form categorical symmetry is a topological, not nec. invertible, $\text{codim} = (p+1)$ operator.

It is called a "p-form" because it assigns charges to p-dim operators:



Lemma: Suppose O, O' are connected by a $(p-1-\dim)$ interface, and Q, Q' are connected by a $(p+2-\text{codim})$ interface. Then $\text{charge}(Q, O) = \text{charge}(Q', O')$.



Completeness of spectrum conjecture:

- If Q, Q' assign the same charges to all p -dim operators, then they are connected by a top'd interface.
- If $\mathcal{O}, \mathcal{O}'$ have the same charges under all p -form categorical symmetries, then they are connected by an interface.

Thm [JF-Reader]:

This is true for topological* quantum field theories.

↪ T_{MN} is central, so all operators are top'd.

To a great extent, GCS = "applied TQFT." E.g.:

"Anomaly Inflow": { QFTs w/ cat'd \mathcal{L} symmetry }


nD \rightarrow { b.c.'s for " \mathcal{L} gauge th'y" } \leftarrow an $n+1D$ TQFT.

Sometimes TQFT methods can prove that a TQFT has no top'd b.c.s. This implies forced gaplessness for QFTs w/ \mathcal{L} symmetry. E.g.:








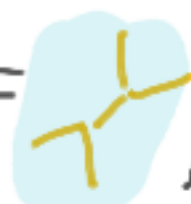


Thm [JF-Reutter]: If a QFT has a framing anomaly \leftarrow an inevitable $n+1D$ forced th'y other than an Atiyah-Kervaire invariant, then it is gapless. and these do have gapped b.c.s.

Finite global categorical symmetries are axiomatized by higher fusion categories. A fusion n -category is a collection of operators and interfaces of $\dim=0, \dots, n$ in $n+1$ ambient dimensions.

They should be:

- Fully topological. E.g.,  is valid.

- Closed under \otimes : if  and  are op's, so is .

- Closed under \oplus and higher ide-potents: If  is s.t.  = , then  :=  exists. Ditto in higher dim, e.g.  if  = ,  = .

- Simple vacuum: Only 0-dim ops are multiples of 1.

- Homotopy-finite: All sets of form {operators}/ {interfaces} are finite.

- Semisimple: For every surface op, its cut of line ops is s.s. i.e. every line defect is a \oplus of simples.

Active research questions

- Understand the "fusion rules" for higher fusion categories
- One source of higher fusion cats is homotopy-finite spaces (= finite higher gps) plus cohomology class plus equivariance data. How much more is there?
- Take constructions from homotopy theory, and implement them on higher fusion categories = "quantum homotopy types".
- Number theory: Can higher fusion cats see beyond \mathbb{Q}^{cyc} ?
- Analysis: Do higher fusion cats come from "higher subfactors"?
- Big Question: Move beyond the finite s.s. world.
- Many more!