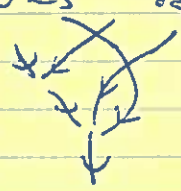


Ideals in derived algebra and boundary conditions in AKSZ-type field theories

Berkeley, Jan 27. Theo Johnson-Freyd

Thank you very much for the invitation. Everything in this talk is dg and over a field of characteristic zero.

I. Probably most of you know what is an operad, but I will review the basic idea. Operads parameterize algebras. The way they do this is as follows. Part of the data of an operad \mathcal{P} is a collection of cochain complexes $\mathcal{P}(n)$, $n \in \mathbb{N}$, of " n -to-1 operators" or "operators of arity n ". Each complex should have an action of symmetric group S_n : if $(a, b) \mapsto a \star b$ is an operation in an algebra, $(a, b) \mapsto b \star a$ should also be. Moreover, there are various composition maps $\mathcal{P}(n) \otimes \mathcal{P}(m) \rightarrow \mathcal{P}(n+m-1)$ because surely $a \star (b \star c)$ is also an operation if \star is. These can enforce algebra axioms like associativity or whatever. The general way to parameterize compositions is in terms of rooted trees



describes some composition

$$\mathcal{P}(2) \otimes \mathcal{P}(3) \rightarrow \mathcal{P}(4)$$

The trees also tell you the "associativity" rules that composition and the S_n action should enjoy.

For this talk I will need to generalize the notion of operad in two simultaneous directions.

(1) A dioperad is like an operad except there are many-to-many operators and compositions are parameterized by (non-rooted) directed trees.

$$\begin{aligned}
 \mathcal{P}(m, n) &= m\text{-to-}n \text{ operations. Various compositions} \\
 \mathcal{P}(m_1, n_1) \otimes \mathcal{P}(m_2, n_2) &\rightarrow \mathcal{P}(m_1 + m_2 - 1, n_1 + n_2 - 1).
 \end{aligned}$$

Example: (a) There is a directed LieB:

generated by operations Υ, λ ,
with relations $\Upsilon = -\phi, \lambda = -\delta$,

$$\Upsilon + \text{cycl}z = 0, \quad \lambda + \text{cycl}z = 0,$$

and $\lambda = \mu + \nu + \xi + \eta$.

(b) There is a directed Frob^{ozo} generated by
 Υ, λ , w/ $\phi = \Upsilon, \lambda = \delta, \Upsilon = \Upsilon, \lambda = \lambda, \lambda = \mu$.

"OCO" stands for "open and co-open". More generally, set Frob^{oco}_{d,d'} to be similar where $\deg(\lambda) = d$ and $\deg(\Upsilon) = d'$. Then

$$H^0(\text{any oriented manifold}) \ni \text{Frob}_{0, \dim}^{\text{ozo}}$$

even if the manifold is open.

(2) An arrowed (di) operad is a 2-colored (di) operad with a distinguished operation

$$\downarrow A$$

$$\uparrow B$$

Here's the reason for the name. An algebra for a (di) operad is a representation of it: a cochain complex V and, for every operad a multilinear map, such that composition in P = composition of multilinear maps.

An algebra for an arrowed directed is automatically an arrow $A \xrightarrow{f} B$.

Examples:

Suppose P is an operad. Set \mathcal{P}

$$P \rightarrow^{st} (m, n; m', n') = \begin{cases} P(m; 1), & (m, m', n') = (m, 0; 1) \\ P(m+n; 1), & m, n; 0, 1 \\ 0 & \text{else.} \end{cases}$$

Then $\mathcal{P} \rightarrow^{st}$ algebras are homomorphisms of P -algebras.

Main ex-ple:

If P is any operad, set

$$P \rightarrow^{st} (m, n; m', n') = \begin{cases} 0 & \text{if } m=n'=0 \\ P(m+n; m'+n') & \text{else.} \end{cases}$$

i.e. you ignore the ~~color~~ color except you disallow all-B-m-to-all-A-out.

III. Quasi-free, etc.

It is technically convenient ^{although not strictly necessary,} to disallow operations with ~~no~~ $m=0$ or $n=0$, and to ~~not~~ only allow identities + the assembly $\&$ when $m+n=m'+n'=1$.
I call such operads oco.

An oco arrowed operad is quasi-free if it is free as a \mathbb{Z} -graded (not nec. dg) object on a generating complex $G = \{G(m, n; m', n'), (m+n)(m'+n') \geq 2\}$.

Standard arguments provide a model category structure on oco objects in which weak equivalences are fibrations and fibrations are surjections. In this model, quasi-free \Rightarrow cofibrant, and every n -cell is a quasi-free resolution.

Suppose P is \mathbb{Z} -free ^{non-graded} on generators G .
 then there is a \mathbb{Z} -free graded \mathbb{Z} -module P°
~~generated by~~

$$G^\circ(m, n; m', n') = \begin{cases} 0 & m=n'=0 \\ G(m+n; m', n') & [m+n'-1] \text{ else.} \end{cases}$$

The differential is given by a sum over internal colorings.

Theorem: $P^\circ \xrightarrow{\sim} P^\circ$ ~~str~~

~~Proof: some identity~~

Let me try to explain P° by way of example.
 Let's set $P = A_\infty$. Recall this is generated by
 operations μ_n , $n \geq 2$, with no symmetry rules.

μ_3 is the associator, ~~str~~...

So, writing $\tilde{\mu} = \mu_n$, we have

$$\partial \mu_3 = \mu_3 - \mu_3, \quad \partial \mu_4 = \mu_4 - \mu_4 - \mu_4 + \mu_4$$

etc.

Then in A_∞° , we have various operations:

$$\begin{array}{c} B \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}, \quad \begin{array}{c} A \\ | \\ B \end{array}$$

$$\begin{array}{c} B \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}, \quad \begin{array}{c} B \\ \diagdown \quad \diagup \\ A \\ \hline A \end{array}, \quad \begin{array}{c} B \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}, \quad \begin{array}{c} B \\ \diagdown \quad \diagup \\ A \\ \hline A \end{array}, \quad \begin{array}{c} A \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}, \quad \begin{array}{c} A \\ \diagdown \quad \diagup \\ A \\ \hline A \end{array}, \quad \begin{array}{c} A \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}$$

$$\partial \begin{array}{c} B \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array} = 0, \quad \partial \begin{array}{c} B \\ \diagdown \quad \diagup \\ A \\ \hline A \end{array} = 0, \quad \partial \begin{array}{c} A \\ \diagdown \quad \diagup \\ A \\ \hline A \end{array} = 0, \quad \text{but}$$

$$\partial \begin{array}{c} B \\ \diagdown \quad \diagup \\ A \\ \hline B \end{array} = \begin{array}{c} B \\ \diagdown \quad \diagup \\ A \\ \hline B \end{array} + \begin{array}{c} B \\ \diagdown \quad \diagup \\ B \\ \hline B \end{array}$$

Look at the all-Bs part. It makes B into an A_{∞} -algebra.

If you look at $A \begin{array}{c} B \dots B \\ \swarrow \searrow \\ A \end{array}$, it makes A into an A_{∞}^B -algebra.
 actually, if you look at $\begin{array}{c} B \dots B \\ \swarrow \searrow \\ A \end{array}$, it makes A into an (A_{∞}^B) -B-bimodule.

$$\partial \left(\begin{array}{c} A \backslash B \\ \hline B \end{array} \right) = \begin{array}{c} A \backslash B \\ \hline B \end{array} \pm \begin{array}{c} A \backslash B \\ \hline A \end{array}.$$

In general, the $\begin{array}{c} B \dots B \\ \swarrow \searrow \\ B \end{array}$ make $f: A \rightarrow B$ into a bimodule map.

But now we have $\begin{array}{c} A \backslash f \\ \hline B \end{array} / A$ and $\begin{array}{c} A \backslash P \\ \hline B \end{array} / A$.

$$\partial \left(\begin{array}{c} A \backslash A \\ \hline A \end{array} \right) = \text{this difference.}$$

If f were an isomorphism, then these would agree. So we make the difference exact.

All together, A_{∞}^{Δ} -algebras are: an A_{∞} -algebra and some type of ideal.

Here's a sharper statement:

Thm: Suppose P is an \mathbb{K} -operad. Then there is an equivalence

$$\text{(homotopy)} \ P^{\Delta}\text{-algebras} \ \simeq \ \text{homotopy } P \rightarrow \text{algebras.}$$

covering

homotopy P^0 -algebras \cong homotopy $P \rightarrow$ algebras

\downarrow forget
{arrows}

relation
of
exact
triangles.

\downarrow forget
{arrows}

$$(A \xrightarrow{f} B) \longmapsto (B \rightarrow \text{cone}(f)).$$

So this really justifies thinking of P^0 -algebras as P -ideals.

Actually, even part of the theorem should surprise you: I'm telling you that if $A \xrightarrow{f} B$ is a P^0 -algebra, then in particular $\text{cone}(f)$ is a P -algebra, at least when P is an operad. This almost holds in the more general doperad case:

Then: If $A \xrightarrow{f} B$ is a P^0 -algebra for P a doperad, then $\text{cone}(f)$ is a

$P \boxtimes \text{Frob}_{0,1}$ algebra.

$P \boxtimes \text{Frob}_{0,1}$ is almost P , just with some different degrees. Here \boxtimes is the Burman-Voght tensor product defined s.t.

$$(P \boxtimes Q)(m, n) = P(m, n) \otimes Q(m, n).$$

Its defining property is that if A is a P -alg and X is a Q -alg, then $A \boxtimes X$ is a $P \boxtimes Q$ -alg.

Example: LieB : $\boxtimes \text{Frob}_{2,0}$ is the version of LieB : here the bracket has degree 2 and the cobracket 2' (up to a sign error).

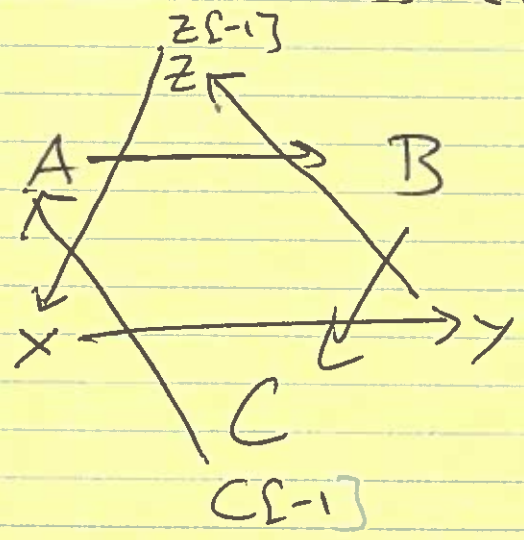
IV. Oh, this reminds me: have you ever wondered how notion of exact triangles interacts with tensor product? See,

~~A~~ $A \rightarrow B$ and $X \rightarrow Y$
 are arrows in ~~A~~
 $A \otimes X \rightarrow B \otimes Y$

is an arrow, but certainly the tensor product of exact triangles is not exact, ~~how~~
~~so it is not so ~~that~~, ~~what~~~~

But if you have two exact triangles
 $A \rightarrow B \rightarrow C$ and $X \rightarrow Y \rightarrow Z$,

you can draw them as an exact Star of David:



and the

Exercise: $\text{Cone}(\begin{smallmatrix} \rightarrow \\ \otimes \\ \rightarrow \end{smallmatrix}) \cong \text{Cone}(\begin{smallmatrix} \downarrow \otimes \\ \downarrow \end{smallmatrix}) \cong \text{Cone}(\begin{smallmatrix} \uparrow \\ \otimes \\ \uparrow \end{smallmatrix})$

This is axiom T3 in May's defn of "symmetric ~~triangulated~~ triangulated category".
~~So there is a ~~relation~~~~

N.B.: One triangle is \odot , the other \ominus , so "forward" nature of one ~~is~~ must be accompanied by "backward" of the other.

Correspondingly you can tensor around depends do get around.

so that if $A \rightarrow B$ is P and $X \rightarrow Y$ is Q ,
~~the~~ $\text{cone}(A \otimes X \rightarrow B \otimes Y)$ is $P \boxtimes Q$.

V. Any world = algebras, operads, ...
 has an operator called the Bar Dual.
 It is defined, up to degree conventions, by

$\mathbb{D}P =$ q-free thing on \underline{P}^* [shift somehow]
 with ∂ encoding composition in P .
 where $\underline{P} = P$ w/o identities.
 For me it's convenient to shift ~~by \mathbb{Z}~~ by \mathbb{Z}

$$P^*(m, n; m', n') [m+n' - \mathbb{Z}].$$

It's called duality because $\mathbb{D}^2 P \xrightarrow{\sim} P$.

You might know the following universal property
 for ~~dis~~ ordinary operads: if P is an operad, $\mathbb{D}P$
~~is universal for~~ is universal for
~~a rep~~ a rep

$$L_{\infty} \rightarrow \mathbb{D}P \otimes P \otimes \mathbb{D}P$$

i.e. if \mathcal{A} is a P alg and \mathcal{B} is a $\mathbb{D}P$ alg,
~~the~~ $\mathcal{A} \otimes \mathcal{B}$ is L_{∞} , and $\mathbb{D}P$ is universal

Thm: For arrowed dioperads, $\mathbb{D}P$ is
 universal ~~st.~~ equipped with a rep

$$L B_{\infty} \rightarrow P \boxtimes \mathbb{D}P$$

$$\parallel$$

$$\mathbb{D} \text{Frob}_{0,1}^{0,0}$$

\parallel
 ∞ -version of Lie bialgebras.

Example: By construction,

$$(\mathbb{D}P)^\vee \cong \mathbb{D}(P^\vee)$$

VI. I will skip quadratic duality. This story is just like in other situations if you know what it means for an operad (or algebra or whatever) to be Koszul, then you know the associated duality result.

Thm: If P is Koszul, so is P^{dual}

~~This makes it easy~~
Koszulity is a tool for computing minimal resolutions.

VII. Relative Poincaré Duality

Now I can get to maybe the central example.

Suppose M is a compact, ^{oriented} d -dim manifold.
This of course $H^*(M)$ is a Frobenius algebra:
you have the cup product, but also a coproduct coming from Poincaré duality.

What if M has boundary? Then you have a rep of commutative algebras

$$\begin{array}{ccc} H^*(M) & & H^*(\partial M) \\ & \searrow & \uparrow \\ & & H^*(M) \end{array}$$

Also you can look at relative cohomology, and you get a rep of coalgebras

$$\begin{array}{ccc} & & H^*(M) \\ & \nearrow & \\ H^*(M, \partial M) & & \end{array}$$

Look at relative cohomology $H^0(m; \partial m)$.

Poincaré duality says $H^0(m; \partial m) \cong H_0(m)$ [shift]
and so $H^0(\partial m) \cong H_0(\partial m)$ [shift]

and in fact

$$H^0(\partial m) \rightarrow H^0(m; \partial m)$$

\exists a map of cochain codes (up to shift).

More is true: In fact,

$$H^0(m; \partial m) \rightarrow H^0(m)$$

\exists a Frobenius Δ -str algebra.

This lifts to (de Rham) cochain level. model

$$H^0(m) \cong \Omega_{\text{dR}}^0(m),$$

$$H^0(m; \partial m) \cong \Omega_{\text{cpt}}^0(m).$$

Defn: A quasi-local many-to-many operator
on de Rham cochains \exists a family of operators ^{low energy limit}
parametrized by an energy scale
~~length~~ energy scale E such that
as $E \rightarrow \infty$ the support (of the integral kernel)
becomes close to the diagonal.

Δ -inequality \Rightarrow these form an (arrows) dispersed

\mathbb{R}^d with the same
reg end on ∂m .

Thm (Cochain-level ^{relative} Poincaré duality).

There is a unique map

$$\{ \text{homotopy } \text{Frob}_{0,2}^{\Delta} \} \rightarrow \{ \alpha\text{-loc operators} \}$$

s.t. the composition

$$\text{homotopy } \text{Frob}_{0,2}^{\Delta} \rightarrow \{ \alpha\text{-loc operators} \} \xrightarrow{\varepsilon=1} \text{End}(\Omega_{\text{cpt}}(M) \rightarrow \Omega(M))$$

$$\xrightarrow{\sim} \text{End}(H(M; \mathbb{Z}) \rightarrow H(M))$$

induces the Frob^{Δ} -str on H° .

VIII - Here's another example. Let me tell you about LB_{∞} -algebras. V . I will ignore degree shifts then V is an L_{∞} -alg and a L_{∞} -coalg in a compatible way.

Take just V as an L_{∞} alg. Then

$$\widehat{\text{Sym}}(V^*[\mathbb{1}]) = \text{CE}^{\circ}(V, \mathbb{1}, \mathbb{J})$$

↑ power series

is an algebra. It is the "functions" on an "infinite-dim" dg manifold w/ coordinate chart V .

V^* is a L_{∞} -alg. The compatibility says:

$\text{CE}^{\circ}(V, \mathbb{1}, \mathbb{J})$ is an L_{∞} -alg,

where the L_{∞} -operators on V^* are extended as multiderivations, i.e.

$\text{CE}^{\circ}(V, \mathbb{1}, \mathbb{J})$ is a Poisson algebra,

i.e. "the infinitesimal manifold V "

is Poisson.

There are some degree shifts.

Suppose Z is Poisson inf manifold aka ~~lie alg~~ ~~LB~~ (homotopy) LB algebra.

What is a coisotropy? Should have

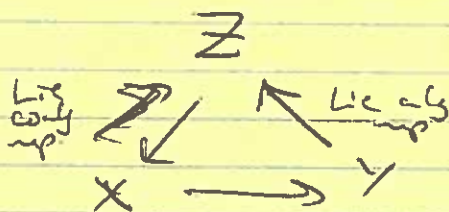
$Y \rightarrow Z$ a "sub" (i.e. map of Lie algs)

and $\ker(O(Z) \rightarrow O(Y)) \hookrightarrow O(Z)$

$\overset{||}{CE(Z)} \rightarrow CE(Y)$

a sub Lie alg. Unpacking:

$Z \rightarrow X = \text{cone}(Y \rightarrow Z)$ should be a map of low-coalgebras.



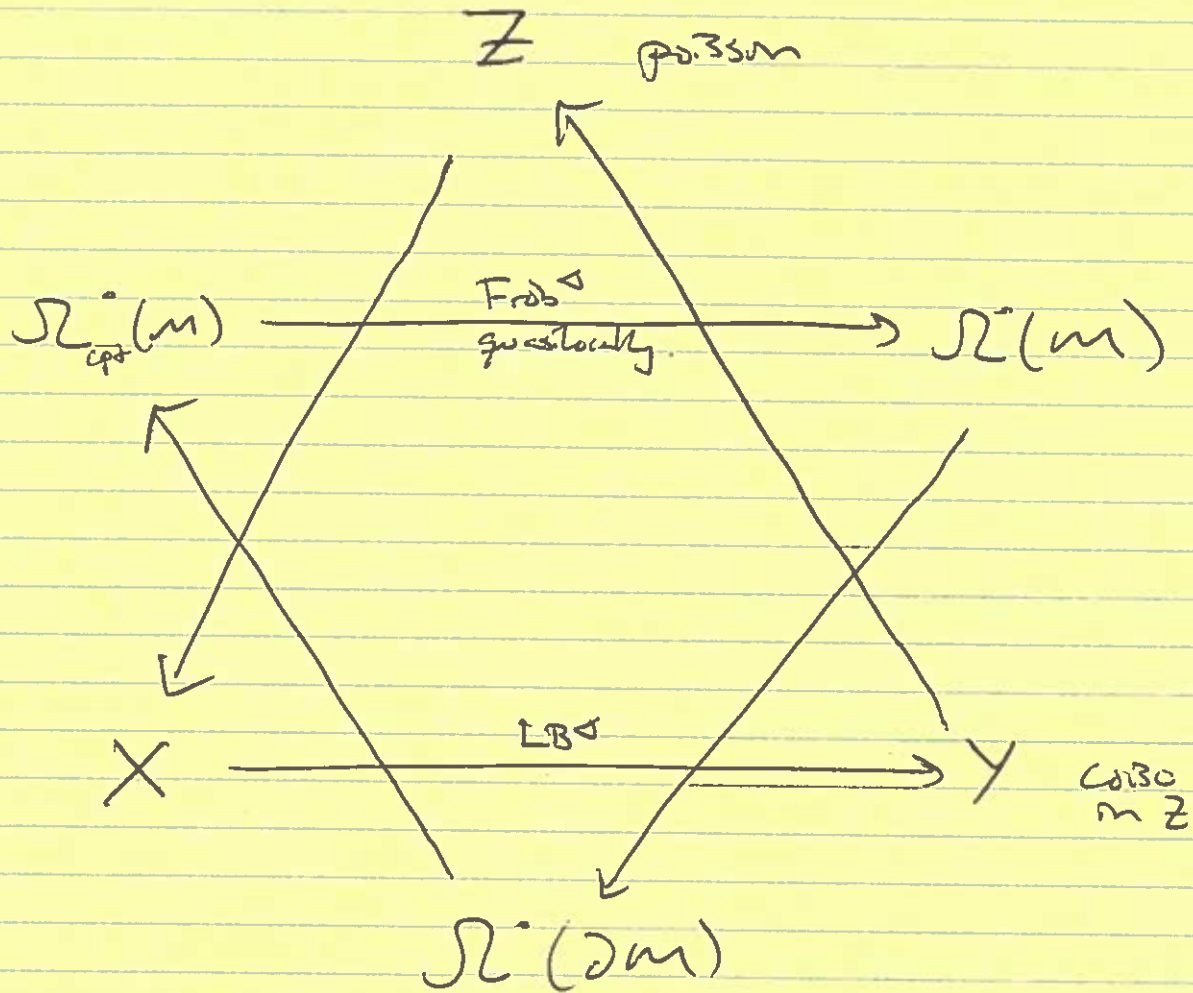
Proposed Defn: ~~Under the data of~~ A map $Y \rightarrow Z$

is coiso in Poisson iff

$\text{cone}(f) \rightarrow Y$ is a LB^0 - (homotopy-) algebra.

~~i.e. the data of~~

IX. Poisson AKSZ with coiso boundary.



Exercise: The invariant V -space — the tensor product of the exact \mathfrak{g} -of-Devi — is also

$$\begin{aligned}
 & \Omega^0(M) \otimes_{\mathbb{Z}} X^h \quad \Omega^0(\partial M) \otimes Y \\
 & \# \quad \Omega^0(\partial M) \otimes_{\mathbb{Z}} \\
 & = \text{maps}(M, \mathbb{Z}) \quad X^h \quad \text{maps}(\partial M, Y) \\
 & \quad \quad \quad \text{maps}(\partial M, \mathbb{Z}) \\
 & = \text{maps} \left(\cdot, \begin{matrix} \circlearrowleft \\ \mathbb{Z} \end{matrix} \right) .
 \end{aligned}$$

Since $LB^0 \cong \mathbb{D} \text{Frob}^0$ (since $LB \cong \mathbb{D} \text{Frob}$),

this invariant space is LB , i.e. ^(shifted) Poisson.

This gives a (shifted) Poisson structure on the space of fields in the Bulk-boundary system (of the expected shift).

The Poisson structure satisfies appropriate locality - condition. (Becomes local in UV limit.)

N.B.: This is classical, not quantum field theory.
 \mathbb{D} -operads \leftrightarrow tree-level Feynman diagrams.